

Multirate GARK schemes for multiphysics problems

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Summary. Multirate GARK schemes define a multirate extension of GARK schemes, generalized additive Runge-Kutta schemes. These allow for exploiting multirate behaviour in both the right-hand sides and in the components in a rather general setting, and are thus especially useful for coupled problems in a multiphysics setting.

We apply MGARK schemes to a benchmark example from thermal-electrical coupling, characterized by a slow and fast part with a stiff and non-stiff characteristic, resp. We test two MGARK schemes: an IMEX method, which makes fully use of the different dynamics and stability properties of the coupled system; and a fully implicit schemes, which inherits the stability properties from both underlying schemes without any coupling constraint.

1 Multirate GARK schemes

We consider a two-way partitioned system

$$y' = f(y) = f^{\{s\}}(y) + f^{\{f\}}, \quad y(t_0) = y_0, \quad (1)$$

with one slow component $\{s\}$, and one active (fast) component $\{f\}$. Note that this setting contains component-wise splitting

$$y = \begin{pmatrix} y_s \\ y_f \end{pmatrix}, \quad f^s = \begin{pmatrix} f_s \\ 0 \end{pmatrix}, \quad f^f = \begin{pmatrix} 0 \\ f_f \end{pmatrix}$$

as a special case.

The slow component is solved with a large step H , and the fast one with small steps $h = H/M$. We will consider the multirate generalization of GARK schemes [3] with M micro steps $h = H/M$, as given in the following

Definition 1 (Multirate GARK method [4]). *One macro-step of a generalized additive multirate Runge-Kutta method with M equal micro-steps reads*

$$\begin{aligned} Y_i^{\{s\}} &= y_n + H \sum_{j=1}^{s\{s\}} a_{i,j}^{\{s,s\}} f^{\{s\}} \left(Y_j^{\{s\}} \right) + \\ &\quad + h \sum_{\lambda=1}^M \sum_{j=1}^{s\{f\}} a_{i,j}^{\{s,f,\lambda\}} f^{\{f\}} \left(Y_j^{\{f,\lambda\}} \right), \\ Y_i^{\{f,\lambda\}} &= y_n + h \sum_{l=1}^{\lambda-1} \sum_{j=1}^{s\{f\}} b_j^{\{f\}} f^{\{f\}} \left(Y_j^{\{f,l\}} \right) + \end{aligned}$$

$$\begin{aligned} &+ H \sum_{j=1}^{s\{s\}} a_{i,j}^{\{f,s,\lambda\}} f^{\{s\}} \left(Y_j^{\{s\}} \right) + \\ &+ h \sum_{j=1}^{s\{f\}} a_{i,j}^{\{f,f\}} f^{\{f\}} \left(Y_j^{\{f,\lambda\}} \right), \\ &\quad \lambda = 1, \dots, M, \\ y_{n+1} &= y_n + h \sum_{\lambda=1}^M \sum_{i=1}^{s\{f\}} b_i^{\{f\}} f^{\{f\}} \left(Y_i^{\{f,\lambda\}} \right) + \\ &+ H \sum_{j=1}^{s\{s\}} b_j^{\{s\}} f^{\{s\}} \left(Y_j^{\{s\}} \right). \end{aligned}$$

The base schemes are Runge-Kutta methods, $(A^{\{f,f\}}, b^{\{f\}})$ for the slow component and $(A^{\{s,s\}}, b^{\{s\}})$ for the fast component. The coefficients $A^{\{s,f,\lambda\}}$, $A^{\{f,s,\lambda\}}$ realize the coupling between the two components.

1.1 Order conditions

The MGARK scheme can be written as a GARK scheme [3] over the macro-step H with the fast stage vectors $Y^{\{f\}} := [Y^{\{f,1\}T}, \dots, Y^{\{f,M\}T}]^T$. The corresponding Butcher tableau reads

$\frac{1}{M}A^{\{f,f\}}$	0	...	0		$A^{\{f,s,1\}}$
$\frac{1}{M}\mathbf{1}b^{\{f\}T}$	$\frac{1}{M}A^{\{f,f\}}$...	0		$A^{\{f,s,2\}}$
\vdots	\vdots	\ddots	\vdots		\vdots
$\frac{1}{M}\mathbf{1}b^{\{f\}T}$	$\frac{1}{M}\mathbf{1}b^{\{f\}T}$...	$\frac{1}{M}A^{\{f,f\}}$		$A^{\{f,s,M\}}$
$\frac{1}{M}A^{\{s,f,1\}}$	$\frac{1}{M}A^{\{s,f,2\}}$...	$\frac{1}{M}A^{\{s,f,M\}}$		$A^{\{s,s\}}$
$\frac{1}{M}b^{\{f\}T}$	$\frac{1}{M}b^{\{f\}T}$...	$\frac{1}{M}b^{\{f\}T}$		$b^{\{s\}T}$

Therefore the order conditions for MGARK schemes can be derived from the corresponding ones for GARK schemes. Up to order two the order conditions given in Table 1 have to be fulfilled.

1.2 Stability

We consider systems (1) where each of the component functions is dispersive (with constants $v^{\{s\}} < 0$, $v^{\{f\}} < 0$):

p	order condition
1	$b^{\{s\}} T \mathbf{1} = 1$ $b^{\{f\}} T \mathbf{1} = 1$
2	$b^{\{s\}} T A^{\{s,s\}} \mathbf{1} = \frac{1}{2}$ $b^{\{s\}} T \left(\sum_{\lambda=1}^M A^{\{s,f,\lambda\}} \right) \mathbf{1} = \frac{M}{2}$ $b^{\{f\}} T A^{\{f,f\}} \mathbf{1} = \frac{1}{2}$ $b^{\{f\}} T \left(\sum_{\lambda=1}^M A^{\{f,s,\lambda\}} \right) \mathbf{1} = \frac{M}{2}$

Table 1. Order conditions for MGARK schemes.

$$\begin{aligned} \langle f^{\{s\}}(y) - f^{\{s\}}(z), y - z \rangle &\leq \mathbf{v}^{\{s\}} \|y - z\|^2, \\ \langle f^{\{f\}}(y) - f^{\{f\}}(z), y - z \rangle &\leq \mathbf{v}^{\{f\}} \|y - z\|^2, \end{aligned}$$

with respect to the same scalar product $\langle \cdot, \cdot \rangle$. As for two solutions $y(t)$ and $\tilde{y}(t)$ of (1), each starting from a different initial condition, the norm of the solution difference $\Delta y(t) = \tilde{y}(t) - y(t)$ is non-increasing, we demand a similar property from the numerical approximations: the MGARK scheme is said to be nonlinearly stable, if the inequality

$$\|y_{n+1} - \tilde{y}_{n+1}\| \leq \|y_n - \tilde{y}_n\|$$

holds for any two numerical approximations y_{n+1} and \tilde{y}_{n+1} obtained by applying the scheme to the ODE (1) with dispersive functions and with initial values y_n and \tilde{y}_n .

As a consequence of stability theory for GARK schemes, an MGARK scheme applied to a component-wise partitioned right-hand side is nonlinearly stable, if both base schemes are algebraically stable.

1.3 Two simple MGARK schemes for multiphysics application

In general, one is interested in a rough approximation of coupled multiphysics problems, which reflect the impact of the couplings of both systems. Hence we restrict to MGARK schemes of order 2:

- **MGARK-IMEX-2:** The implicit-explicit version solves the fast, stiff part with an implicit base scheme, and the slow, non-stiff part with an explicit one. The coefficients are given by

$$\begin{aligned} b^{\{s\}} &= \begin{pmatrix} \frac{1}{2} & \frac{1}{2} \end{pmatrix}, \quad A^{\{s,s\}} = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}, \\ A^{\{s,f,1\}} &= \begin{pmatrix} 0 \\ M \end{pmatrix}, \\ A^{\{s,f,\lambda\}} &= \begin{pmatrix} 0 \\ 0 \end{pmatrix} \quad \forall \lambda = 2, \dots, M, \end{aligned}$$

$$b^{\{f\}} = 1, \quad A^{\{f,f\}} = \frac{1}{2},$$

$$A^{\{f,s,\lambda\}} = \begin{pmatrix} \frac{1}{2} & 0 \end{pmatrix} \quad \forall \lambda = 1, \dots, M.$$

Note that only the fast part is algebraically stable, but neither the slow part and the joint system.

- **MGARK-IMIM-2:** To get an overall stable scheme, both parts are solved by an implicit base scheme. The coefficients are given by

$$b^{\{s\}} = \begin{pmatrix} 0 & 1 \end{pmatrix}, \quad A^{\{s,s\}} = \begin{pmatrix} 0 & 0 \\ 0 & \frac{1}{2} \end{pmatrix},$$

$$A^{\{s,f,1\}} = \begin{pmatrix} 0 \\ \frac{M}{2} \end{pmatrix},$$

$$A^{\{s,f,\lambda\}} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \quad \forall \lambda = 2, \dots, M,$$

$$b^{\{f\}} = 1, \quad A^{\{f,f\}} = \frac{1}{2},$$

$$A^{\{f,s,\lambda\}} = \begin{pmatrix} \frac{1}{2} & 0 \end{pmatrix} \quad \forall \lambda = 1, \dots, M.$$

As both base schemes are algebraically stable, the MGARK method inherits this property for a component-wise partitioning.

2 Benchmark example

We will test both MGARK implementations for the electrical-thermal multiphysics system introduced in [1] with the specifications discussed in [2]. The thermal component defines the slow (and non-stiff) part, the electrical component the fast (and stiff) part of the system.

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